

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems
Quantum
Entanglement

Results

General
Conclusions

Quantifying quantum coherence in molecular magnetic systems

Clebson dos Santos Cruz

Grupo de Informação Quântica e Física Estatística, Universidade Federal
do Oeste da Bahia

Barreiras

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List of Contents

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- 1 Introduction
- 2 Low-Dimensional Molecular Magnetism
 - Hamiltonians and Interactions
 - Thermodynamic Quantities
- 3 Quantum Coherence in Molecular Magnetic Systems
- 4 Results
- 5 General Conclusions

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Introduction

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
 - Technological breakthrough;
 - Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
 - Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
 - Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices;
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional Molecular Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum Coherence in Molecular Magnetic Systems

Quantum
Entanglement

Results

General Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Miniaturization of electronic devices:
- Technological breakthrough;
- Quantum Computation:
 - Outperform classical computers;
 - Exploiting quantum properties
- Demand for the development of novel materials
 - Design and synthesis;
 - Characterization of its quantum properties;
- Molecular magnetic systems \Rightarrow new quantum technologies.

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

NATURE | VOL 410 | 12 APRIL 2001 | www.nature.com

Quantum computing in molecular magnets

Michael N. Leuenberger & Daniel Loss

*Department of Physics and Astronomy, University of Basel, Klingelbergstrasse 82,
4056 Basel, Switzerland*

Shor and Grover demonstrated that a quantum computer can outperform any classical computer in factoring numbers¹ and in searching a database² by exploiting the parallelism of quantum mechanics. Whereas Shor's algorithm requires both superposition and entanglement of a many-particle system³, the superposition of single-particle quantum states is sufficient for Grover's algorithm⁴. Recently, the latter has been successfully implemented⁵ using Rydberg atoms. Here we propose an implementation of Grover's algorithm that uses molecular magnets⁶⁻¹⁰, which are solid-state systems with a large spin; their spin eigenstates make them natural candidates for single-particle systems. We show theoretically that molecular magnets can be used to build dense and efficient memory devices based on the Grover algorithm. In particular, one single crystal can serve as a storage unit of a dynamic random access memory device. Fast electron spin resonance pulses can be used to decode and read out stored numbers of up to 10^5 , with access times as short as 10^{-10} seconds. We show that our proposal should be feasible using the molecular magnets Fe₈ and Mn₁₂.

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

NATURE | VOL 410 | 12 APRIL 2001 | www.nature.com

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- Stability of its quantum properties;
- Decoherence problem;
- Control the quantum properties;
- Novel materials for quantum information processing.

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

NATURE | VOL 410 | 12 APRIL 2001 | www.nature.com

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- Stability of its quantum properties;
- Decoherence problem;
- Control the quantum properties;
- Novel materials for quantum information processing.

Clebson Cruz

Introduction

Low-Dimensional Molecular Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum Coherence in Molecular Magnetic Systems

Quantum
Entanglement

Results

General Conclusions

NATURE | VOL 410 | 12 APRIL 2001 | www.nature.com

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- Stability of its quantum properties;
- Decoherence problem;
- Control the quantum properties;
- Novel materials for quantum information processing.

Clebson Cruz

Introduction

Low-Dimensional Molecular Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum Coherence in Molecular Magnetic Systems

Quantum
Entanglement

Results

General Conclusions

NATURE | VOL 410 | 12 APRIL 2001 | www.nature.com

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- Stability of its quantum properties;
- Decoherence problem;
- Control the quantum properties;
- Novel materials for quantum information processing.

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Low-Dimensional Molecular Magnetism

Introduction

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

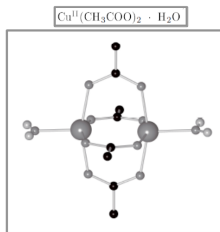
Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

• Brief History



Atomic structure of copper acetate dimer¹

¹BLEANEY, B.; BOWERS, K. D.; *Anomalous paramagnetism of copper acetate*. Proc. R. Soc. A, 212(1119): 451-465, 1952.

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History

- Molecular magnetism

- Chemistry-Magnetism;

- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History

- Molecular magnetism

- Chemistry-Magnetism;
- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
- Molecular magnetism

- Chemistry-Magnetism;
- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
- Molecular magnetism

- Chemistry-Magnetism;
- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
- Molecular magnetism

- Chemistry-Magnetism;
- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
- Molecular magnetism
- Chemistry-Magnetism;
- Non-cooperative;
- Short range interaction;
- Located;
- Orbitals d , f , p ;
- Low Dimensional;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Main Properties and Fundamental Aspects

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

- Brief History
 - Molecular magnetism
 - Fundamental Aspects
- Easy Synthesis;
 - Versatility;
 - Biocompatibility;
 - Electrical insulation;
 - Solubility;
 - Quantum effects;

Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

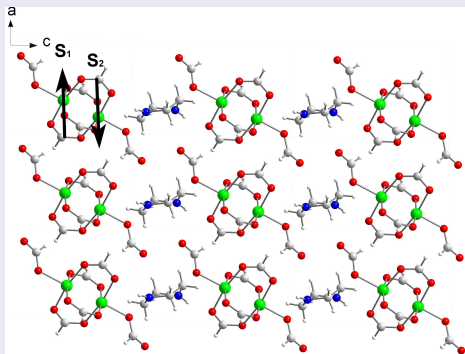
Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Geometric Configurations



Crystal structure of $\text{Cu}_2 (\text{HCOO})_4 (\text{HCOOH})_2$
 $(\text{C}_4\text{H}_{10}\text{N}_2)$

Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Hamiltonians and Interactions:

- Heisenberg Hamiltonian:

$$\mathbf{H}_{hei} = -J\vec{S}_1 \cdot \vec{S}_2;$$

$$J = E_{AP} - E_P$$

- $J=0$, there is no interaction ($E_{AP} = E_P$).
- $J>0$, parallel alignment ($E_{AP} > E_P$)
- $J<0$, antiparallel alignment ($E_{AP} < E_P$)
- Zeeman:

$$\mathbf{H}_z = \mu_B \vec{B} \cdot \mathbf{g} \cdot \vec{S};$$

Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Hamiltonians and Interactions:

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Hamiltonians and Interactions:

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Hamiltonians and Interactions:

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermodynamic Quantities:

$$\mathcal{H} = -J\vec{S}_1 \cdot \vec{S}_2 - g\mu_B \vec{B} \cdot (\vec{S}_1 + \vec{S}_2)$$

$$\mathcal{E}_{s,m_s} = -\frac{1}{2}J[s(s+1) - s_1(s_1+1) - s_2(s_2+1)] - g\mu_B B m_s$$

where $|s_1 - s_2| \leq s \leq s_1 + s_2$; and $m_s = -s, \dots, +s$.

$$Z(T, B) = \sum_i e^{-\beta \mathcal{E}_i}$$

$$F(T, B) = -k_B T \ln(Z(B, T))$$

$$M(T, B) = -\frac{\partial}{\partial B} F(T, B)$$

$$\chi_{ij} = \lim_{B \rightarrow 0} \mu_0 \frac{dM_j}{dB_i} = \frac{2N(g\mu_B)^2}{k_B T} \frac{1}{3 + e^{-J/k_B T}}$$

Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum

Coherence in

Molecular

Magnetic Systems

Quantum

Entanglement

Results

General

Conclusions

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$$\mathcal{H} = -J\vec{S}_1 \cdot \vec{S}_2 - g\mu_B \vec{B} \cdot (\vec{S}_1 + \vec{S}_2)$$

$$\mathcal{E}_{s,m_s} = -\frac{1}{2}J[s(s+1) - s_1(s_1+1) - s_2(s_2+1)] - g\mu_B B m_s$$

where $|s_1 - s_2| \leq s \leq s_1 + s_2$; and $m_s = -s, \dots, +s$.

$$Z(T, B) = \sum_i e^{-\beta \mathcal{E}_i}$$

$$F(T, B) = -k_B T \ln(Z(B, T))$$

$$M(T, B) = -\frac{\partial}{\partial B} F(T, B)$$

$$\chi_{ij} = \lim_{B \rightarrow 0} \mu_0 \frac{dM_j}{dB_i} = \frac{2N(g\mu_B)^2}{k_B T} \frac{1}{3 + e^{-J/k_B T}}$$

Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum

Coherence in

Molecular

Magnetic Systems

Quantum

Entanglement

Results

General

Conclusions

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum

Coherence in

Molecular

Magnetic Systems

Quantum

Entanglement

Results

General

Conclusions

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Low-Dimensional Molecular Magnetism

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

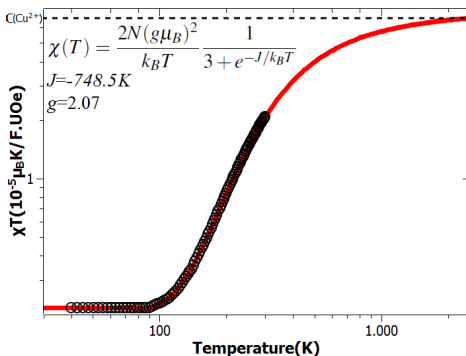
Results

General
Conclusions

Experimental Data

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Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence in Molecular Magnetic Systems

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

- Remarkable resource for quantum information processing;
- Let us consider the Hilbert space of a composite system:

$$\mathbb{H} = \mathbb{H}_1 \otimes \mathbb{H}_2 \otimes \cdots \otimes \mathbb{H}_n$$

- An entangled state can be defined as:

$$\begin{aligned} |\Psi\rangle &\neq |\phi_1\rangle \otimes |\phi_2\rangle \otimes \cdots \otimes |\phi_n\rangle \\ \rho &= \sum_i p_i \rho_1^{(i)} \otimes \rho_2^{(i)} \otimes \cdots \otimes \rho_n^{(i)} \end{aligned}$$

- The best knowledge of a whole does not include the knowledge of its parts.

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

- Let us consider the following example:

Event	Box A	Box B	Probability
1	0	1	50%
2	1	0	50%

$$|\phi_A\rangle = |\phi_B\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|0\rangle_A \otimes |1\rangle_B + |1\rangle_A \otimes |0\rangle_B) \neq |\phi_A\rangle \otimes |\phi_B\rangle .$$

- Entangled states cannot be simulated or represented from classical correlations;
- Entangled \Rightarrow Quantum Coherence + Classical Correlations;

Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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- Entangled \Rightarrow Quantum Coherence + Classical Correlations;

Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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- Entangled \Rightarrow Quantum Coherence + Classical Correlations;

Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Correlations in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

- Let us consider a Heisenberg dimer:

$$\mathbf{H} = -J\vec{S}_A\vec{S}_B$$

$$\rho_{AB} = e^{-\beta\mathbf{H}}/Z$$

$$\rho_{AB}(T) = \frac{1}{4} \begin{bmatrix} 1 + c(T) & & & \\ & 1 - c(T) & 2c(T) & \\ & 2c(T) & 1 - c(T) & \\ & & & 1 + c(T) \end{bmatrix}$$

$$c(T) = \langle \vec{S}_A^{(\alpha)} \otimes \vec{S}_B^{(\alpha)} \rangle = \frac{2k_B T}{N(g\mu_B)^2} \chi(T) - 1$$

- SPOILER ALERT:** Coherence \Rightarrow All off-diagonal elements!

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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$$|\Psi\rangle \neq |\phi_1\rangle \otimes |\phi_2\rangle$$

$$\rho \neq \sum_i p_i \rho_A^{(i)} \otimes \rho_B^{(i)} \otimes \rho_C^{(i)} \otimes \dots$$

$$\mathbf{H}_{hei} = -J \vec{S}_1 \cdot \vec{S}_2; \quad E_{S_{AP}} - E_{S_P} = J$$

$J > 0$, parallel alignment ($E_{S_{AP}} > E_{S_P}$)

$J < 0$, antiparallel alignment ($E_{S_{AP}} < E_{S_P}$)

E_{s, m_s}	$ s, m_s\rangle$	$\{ m_{s1}, m_{s2}\rangle\}$	
0	$ 1, 1\rangle$	$ \uparrow\uparrow\rangle$	separable
0	$ 1, -1\rangle$	$ \downarrow\downarrow\rangle$	separable
0	$ 1, 0\rangle$	$ \uparrow\downarrow\rangle + \downarrow\uparrow\rangle$	entangled
J	$ 0, 0\rangle$	$ \uparrow\downarrow\rangle - \downarrow\uparrow\rangle$	entangled

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Thermal Entanglement

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence:

- Entangled states exhibit quantum and classical correlation,
- Does not encompass all quantum correlations.
- Quantum coherence, arising from quantum superposition.
- Coherence has an important role in many quantum information processes.

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence:

- Entangled states exhibit quantum and classical correlation,
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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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- Coherence has an important role in many quantum information processes.

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence:

- Geometric approaches are widely used to characterize and quantify the quantum correlations.

- *Phys. Rev. Lett.* 113, 140401 (2014):

From the minimal distance $D(\rho, \sigma)$, between the quantum state ρ and a set $\{\sigma = \sum_k^d |k\rangle\langle k| \in \mathcal{I}\}$ of incoherent state.

$$\mathcal{C}_D = \min_{\{\sigma \in \mathcal{I}\}} D(\rho, \sigma)$$

the l_1 trace norm can be a reliable measurement of quantum coherence as

$$\mathcal{C}_{l_1} = \min_{\sigma \in \mathcal{I}} \|\rho - \sigma\|_{l_1} = \sum_{i \neq j} |\langle i | \rho | j \rangle| .$$

Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Quantum Coherence in Molecular Magnetic Systems

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Results

Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

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Introduction

Low-Dimensional
Molecular
Magnetism

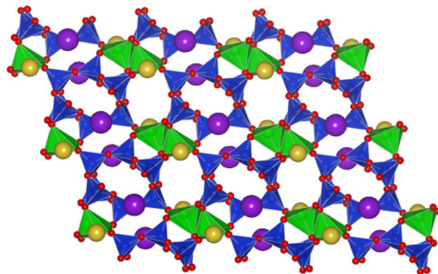
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Crystal structure of the Metal Silicate Framework



Journal of Solid State Chemistry, 182(2), 253-258.

Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

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Introduction

Low-Dimensional
Molecular
Magnetism

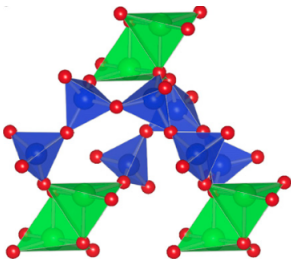
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin 1/2 dimer;
- 2 qubit system;
- Dimers are magnetically isolated;
- Separated by two SiO_4 corners;

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Introduction

Low-Dimensional
Molecular
Magnetism

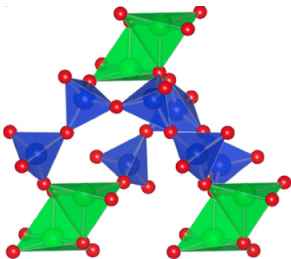
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin 1/2 dimer;
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Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

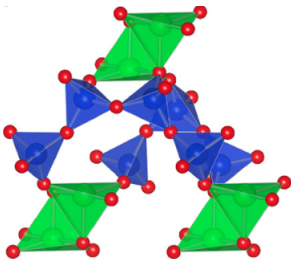
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin 1/2 dimer;
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Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

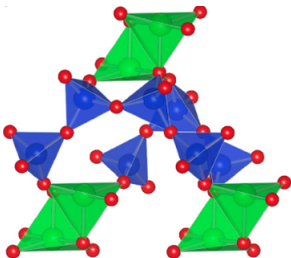
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin 1/2 dimer;
- 2 qubit system;
- Dimers are magnetically isolated;
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Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

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Introduction

Low-Dimensional
Molecular
Magnetism

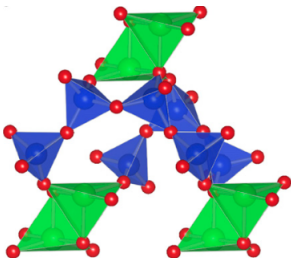
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin $1/2$ dimer;
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- Dimers are magnetically isolated;
- Separated by two SiO_4 corners;

Metal Silicate Framework $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

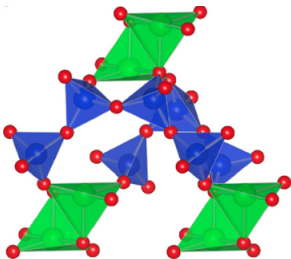
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Journal of Solid State Chemistry, 182(2), 253-258.

- Cu^{2+} ions - d^9 electronic configuration;
- a Heisenberg spin $1/2$ dimer;
- 2 qubit system;
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Quantum Coherence of $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence

- Let us consider a Heisenberg dimer:

$$\rho_{AB} = e^{-\beta \mathbf{H}} / Z \text{ with } \mathbf{H} = -J \vec{S}_A \vec{S}_B$$

$$\rho_{AB}(T) = \frac{1}{4} \begin{bmatrix} 1 + c(T) & & & \\ & 1 - c(T) & 2c(T) & \\ & 2c(T) & 1 - c(T) & \\ & & & 1 + c(T) \end{bmatrix}$$

$$c(T) = \frac{2k_B T}{N(g\mu_B)^2} \chi(T) - 1$$

$$C(T) = \left| \frac{2k_B T \chi(T)}{N_A g^2 \mu_B^2} - 1 \right|$$

Quantum Coherence of $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Quantum Coherence

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Quantum Coherence of $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Quantum Coherence of $\text{KNaCuSi}_4\text{O}_{10}$

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

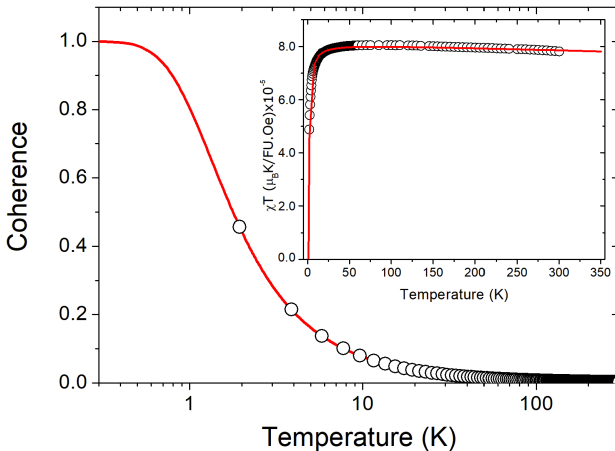
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Acknowledgements: A.M. dos Santos - Oak Ridge National Laboratory, Tennessee USA

Influence of the External Pressure

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems
Quantum
Entanglement

Results

General
Conclusions

First-Principles Calculations

- Kohn-Sham equations were solved;

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V_s(\vec{r}) \right] \phi_i(\vec{r}) = E_i \phi_i(\vec{r})$$

$$\text{where } n(\vec{r}) = \sum_i^N |\phi_i(\vec{r})|^2$$

$$V_s(\vec{r}) = V(\vec{r}) + \int \frac{e^2 n_s(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3 r' + V_{xc}[n_s(\vec{r})]$$

- Crystal structure was optimized;
- Third-order Birch-Murnaghan equation;

Influence of the External Pressure

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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Influence of the External Pressure

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

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- Third-order Birch-Murnaghan equation;

Influence of the External Pressure

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

First-Principles Calculations

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Influence of the External Pressure

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

First-Principles Calculations

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First-Principles Calculations

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

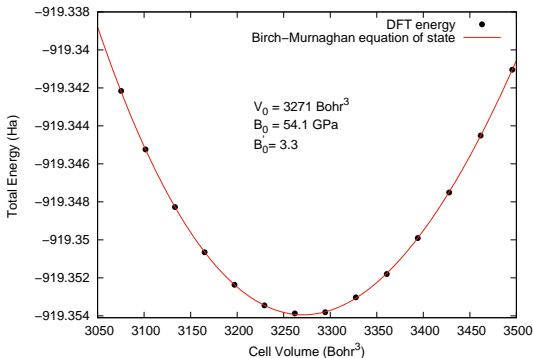
Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

$$E(V) = E_0 + \frac{9V_0B_0}{16} \left\{ \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - 1 \right]^3 B'_0 + \left[\left(\frac{V_0}{V} \right)^{\frac{2}{3}} - 1 \right]^2 \left[6 - 4 \left(\frac{V_0}{V} \right)^{\frac{2}{3}} \right] \right\}$$



Results

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$$J = E_S - E_T$$

Introduction

Low-Dimensional
Molecular
Magnetism

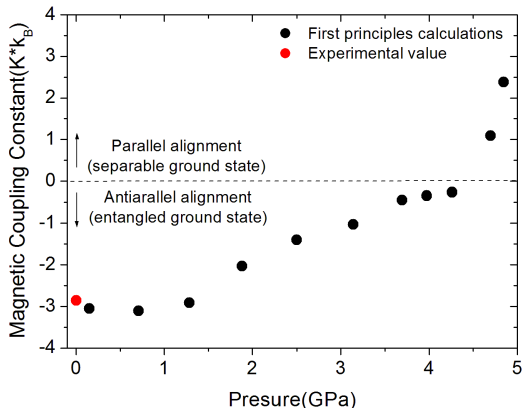
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Results

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

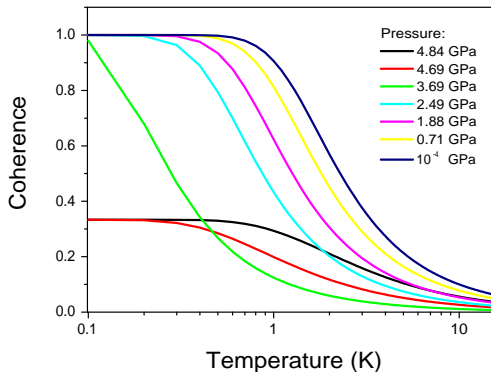
Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Temperature dependence of the quantum coherence calculated for different values of hydrostatic pressure.

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Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Influence of the longitudinal and transverse magnetic field

- Quantum coherence is basis dependent.
- Basis choice depends on the physical problem under investigation.
- Molecular magnetic systems: spin eigenbasis in a certain direction, $\{S_x, S_y, S_z\}$, within a quantum metrology setting.

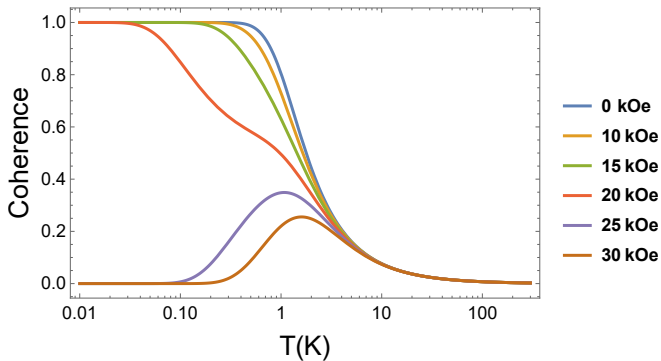
The Hamiltonian that rule this system interacting with an external magnetic field is given by:

$$\mathcal{H} = -J\vec{S}_1 \cdot \vec{S}_2 - \mu_B g \vec{B} \cdot (\vec{S}_1 + \vec{S}_1).$$

Longitudinal Magnetic Field

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$$C_z(T, B_z) = \left| \frac{1 - e^{-4x}}{1 + e^{-4x} + 2 \cosh(\beta h_z)} \right|$$



Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Longitudinal Magnetic Field

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Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

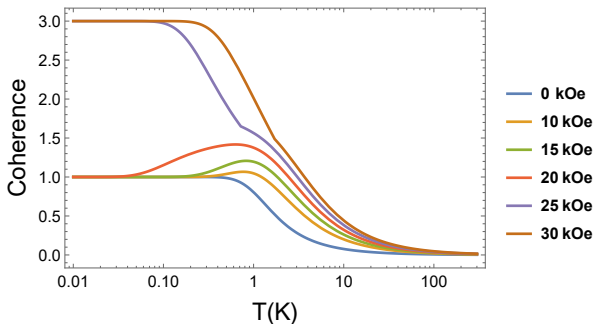
Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

$$C_x(T, B_x) = \frac{e^x}{Z} (|\cosh(\beta h_z) - 1| + 4|\sinh(\beta h_z)| + |\cosh(\beta h_z) - e^{-4x}|)$$



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Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

General Conclusions and Future Works

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Conclusions

- Detection and quantification of quantum coherence in molecular magnets;
- Coherent quantum states at room temperatures;
- Dependence under hydrostatic pressure;
- Control and management of quantum properties;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Conclusions

- Detection and quantification of quantum coherence in molecular magnets;
- Coherent quantum states at room temperatures;
 - Dependence under hydrostatic pressure;
 - Control and management of quantum properties;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Conclusions

- Detection and quantification of quantum coherence in molecular magnets;
- Coherent quantum states at room temperatures;
- Dependence under hydrostatic pressure;
- Control and management of quantum properties;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Conclusions

- Detection and quantification of quantum coherence in molecular magnets;
- Coherent quantum states at room temperatures;
- Dependence under hydrostatic pressure;
- Control and management of quantum properties;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Conclusions

- Detection and quantification of quantum coherence in molecular magnets;
- Coherent quantum states at room temperatures;
- Dependence under hydrostatic pressure;
- Control and management of quantum properties;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Future Works

- Study the creation of quantum logic gates with the molecular magnets;
- Obtain new topologies of molecular magnetic systems with optimized quantum properties;
- Study thermal effects on the quantum coherence;
- Molecular engineering for quantum computing;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Future Works

- Study the creation of quantum logic gates with the molecular magnets;
- Obtain new topologies of molecular magnetic systems with optimized quantum properties;
- Study thermal effects on the quantum coherence;
- Molecular engineering for quantum computing;

General Conclusions and Future Works

Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations

Hamiltonians and
Interactions

Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Future Works

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Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

Future Works

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Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions



Clebson Cruz

Introduction

Low-Dimensional
Molecular
Magnetism

Theoretical
Foundations
Hamiltonians and
Interactions
Thermodynamic
Quantities

Quantum
Coherence in
Molecular
Magnetic Systems

Quantum
Entanglement

Results

General
Conclusions

